

### QUESTIONS

**ING** (interneuronal network gamma) and (pyramidal-interneuronal PING network gamma) are well-established theories of the mechanistic generation of  $\gamma$ -rhythms ( $\approx$  35-100 Hz). (2) However, the capacity of these mechanisms to dynamically respond to external inputs and phase-lock with upstream rhythms has not been studied.

We pose the questions:

- What range of dynamics is possible for these models under periodic forcing?
- How do the properties of these models compare to those of other commonly studied forced oscillators?
- Are these mechanisms well-suited to rapidly establish reliable phase-locked relationships with upstream  $\gamma$ -rhythms?

We address these questions by mathematically analyzing simple ING and PING models and comparing them to **phase oscillators** and **re**laxation oscillators.

#### PHASE OSCILLATOR LIMITATIONS

The **phase oscillator** is the generic form of a 1D periodicallyforced oscillator. By changes of variables, any stable limit cycle under sufficiently weak forcing ( $\epsilon < \epsilon^*$  for some  $\epsilon^* > 0$ ) may be written in this form, as may the LIF and QIF neuron models.

$$\dot{\phi} = 1 + g(\phi)(b + \epsilon I(\Phi))$$
$$\dot{\Phi} = 1$$

Restriction to one dimension imposes a tradeoff between robustness of natural period T to changes in tonic drive b and capacity to lock to weak inputs at a wide range of frequencies. At a given forcing strength  $\epsilon$ , the width  $W^{\epsilon}$  of the interval of forcing periods  $T_I$  which may evoke stable phase locking is bounded by

 $W_{\epsilon} < \sqrt{-\frac{\partial T(0,b,\Phi_0,T_I)}{\partial b}}|_{b=0} \sqrt{\epsilon \sup_{\Phi_0 \in [0,T_I]} \int_0^{T_{0,0}} I(\Phi_0 + t \mod T_I) dt}$ 

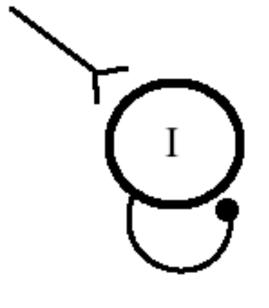
# **GAMMA** ( $\gamma$ ) **RHYTHMS UNDER PERIODIC FORCING**

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## GAMMA-GENERATING MECHANISMS

**ING** is modeled as a synchronous population of inhibitory theta-neurons with phase  $\theta_i$  and tonic excitation  $b_i$ , inhibiting itself with slowly decaying inhibition  $s_i$  and receiving a periodic excitatory input  $\epsilon I(\Phi)$  with period  $T_I$ .

**PING** is modeled like ING, with the addition of quickly-decaying excitation  $s_e$  from synchronous population of excitatory cells with phase  $\theta_e$ , which triggers the I-population. The E-population receives the forcing  $\epsilon I(\Phi)$ , and is inhibited by the I-population.

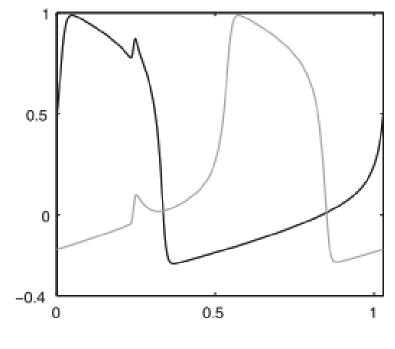


**I-population phase:**  $\tau_i \dot{\theta} = 1 - \cos(\theta_i) + (1 + \cos(\theta_i))G_i$  $G_i = b_i - g_{ii}s_i + \epsilon I(\Phi), b_i > 0$ 

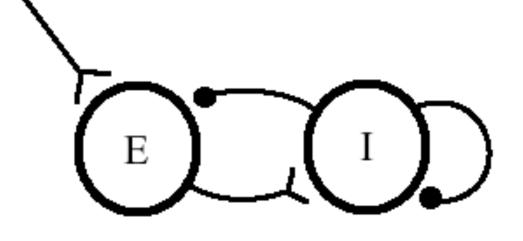
> Inhibition:  $\dot{s_i} = -s_i/\tau_{s_i}$ When  $\theta = \pi$ ,  $s_i$  resets to  $s_i = c(s_i - 1) + 1$

Forcing phase:  $\Phi = 1$ 

A separation of timescales can circumvent this problem by making the  $\epsilon^*$  required for the phase reduction arbitrarily small. Thus, if the ING oscillator has  $\tau_{s_i} \gg \tau$ , it both maintains a robust natural period *T* and phaselocks to inputs at a range of forcing periods unbounded by forcing current and sensitivity.



Regimes of 1:1-2:1 bistability and period-doubing under stronger forcing. (**Right:** period-doubling bifurcation as forcing strength increases, in red)



**E/I-population phases:**  $\tau_{e/i} \dot{\theta}_{e/i} = 1 - \cos(\theta_{e/i}) + (1 + \cos(\theta_{e/i}))G_{e/i}$  $G_e = b_e - g_{ie}s_i + \epsilon I(\Phi), b_e > 0$  $G_i = b_i - g_{ii}s_i + g_{ei}s_e, b_i < 0$ 

> **Excitation/Inhibition:**  $\dot{s}_{e/i} = -s_{e/i}/\tau_{s_i}$ When  $\theta_{e/i} = \pi$ ,  $s_{e/i}$  resets to 1.

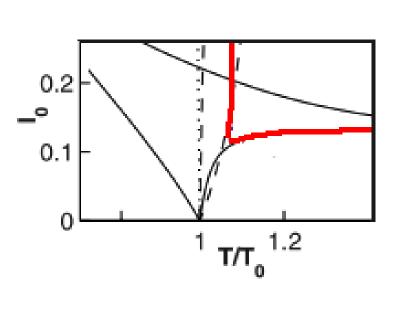
> > Forcing phase:  $\dot{\Phi} = 1$

# **RELAXATION LIMITATIONS**

The separate timescales of the Fitzhugh-Nagumo **relaxation oscillator** give it a robust period while phase-locking robustly to inputs.

Two distinguishing characteristic properties of the forced phase oscillator are:

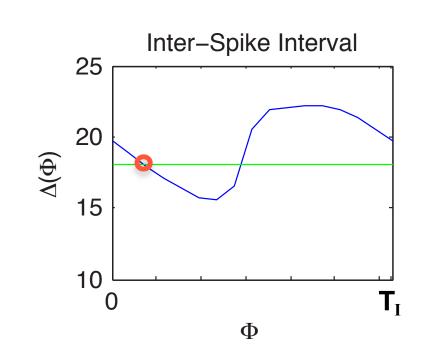
> Regimes of 1:1 bistability under weak forcing. (Left: A bistable pair of 1:1 forced orbits)



Figures and results from (1).

 $( au_{s_i} \gg au).$ 

ble.



**2.** If *c* is small and/or sufficient time is spent under inhibition each cycle, the system attracts to an invariant torus on which perioddoubling and 1:1-2:1 bistability are not possible.

**Right:** Given the conditions above, an invariant torus persists under strong forcing. The same applies to the **PING** model if  $\tau_i$  is small.

# CONCLUSION

The **ING** and **PING** mechanisms are **ideally suited** for responding to upstream  $\gamma$ -rhythms of varying amplitude and frequency by rapidly establishing a **reliable phase-locked relationship**.

REFER	
1.	H. Cro 1539-37
2.	М. а. <i>psychop</i> 2000).



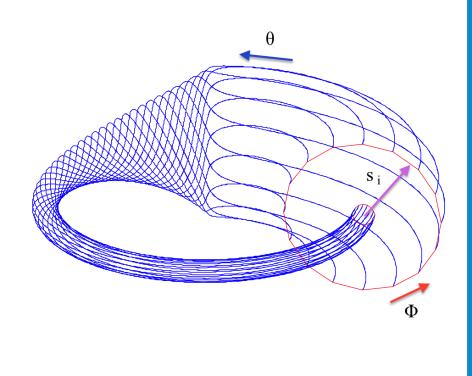
### **ING/PING PROPERTIES**

Like the relaxation oscillator (below), the  $\gamma$ -generating networks maintain a robust natural period while robustly phase-locking to forcing due to separation of time scales

We use the variational equations and return maps for our models (left) to prove two results differentiating them from ordinary relaxation:

Only one phase offset between the ING oscillator and pulsatile forcing is sta-We prove this for square pulses.

> Left: Phase-locking only occurs where the ISI function crosses the forcing period downwards, which in this case can occur only once.



### ENCES

Disier et al., Physical Review E 79, 57–59, ISSN: 3755 (Jan. 2009).

Whittington et al., International journal of physiology 38, 315–36, ISSN: 0167-8760 (Dec.